

Lecture 11 Atomic Spectra

SCPY152, Second Semester 2021-22

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Topics

1. Principle hydrogen spectrum
2. Atom in external fields
3. Fine structure

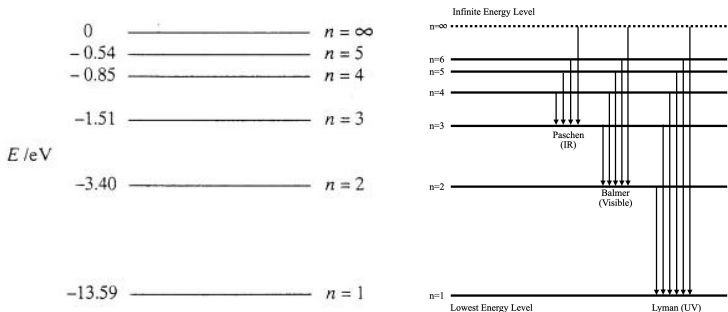
Principle spectrum

Hydrogen energy levels (from Bohr and Schrodinger)

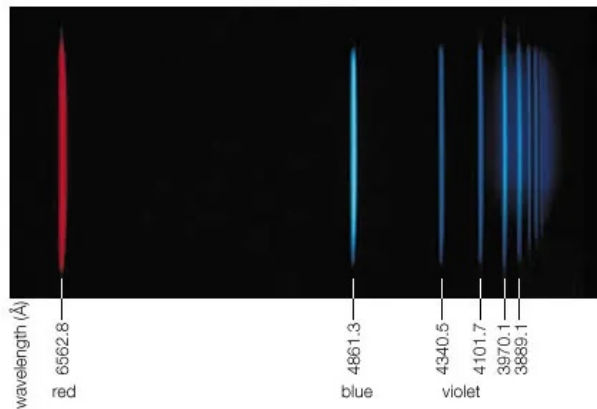
$$E_n = -\frac{1}{2} \frac{Ke^2}{a_0 n^2} = -\frac{1}{2} m_e c^2 \alpha^2 \frac{1}{n^2} = -\frac{13.6\text{eV}}{n^2} \quad (1)$$

with $a_0 = \frac{\hbar^2 c^2}{m_e c^2 Ke^2}$ and $\alpha = \frac{Ke^2}{\hbar c}$

Energy levels appear on the left, and emission transitions appear on the right



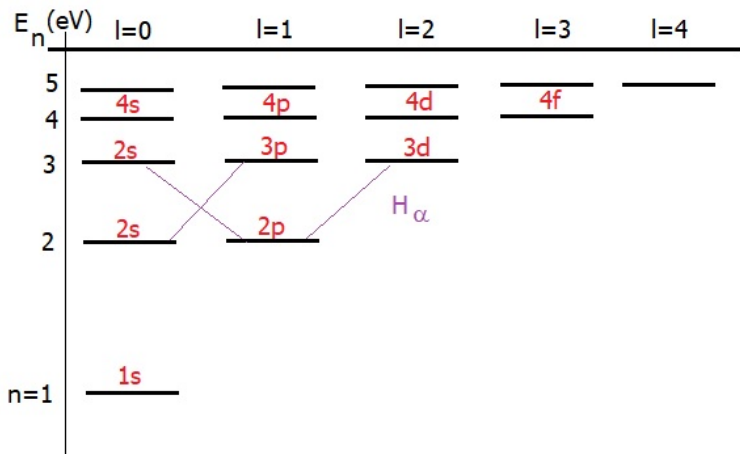
Appearance spectrum in Balmer series



Atomic transition and selection rules

Selection rules of atomic transitions (emission/absorption)

$$\Delta n = \text{any}, \quad \Delta l = \pm 1$$



Hydrogen atom in (weak) external fields

Atom in uniform electric field $\vec{E} = E_0 \hat{z}$: Stark effect

Interaction potential energy of electron (with electric charge $-e$)

$$U = e\vec{E} \cdot \vec{r} = eE_0 r \cos \theta \quad (2)$$

Effect to the s-orbital, i.e., 1s-orbital

$$\langle U \rangle_{100} = \frac{1}{\pi a_0^3} eE_0 \int_0^\infty r^3 e^{-2r/a_0} dr \int_{-1}^1 \cos \theta d \cos \theta \int_0^{2\pi} \theta d\phi = 0 \quad (3)$$

No effect on s-orbital. Effect on p-orbitals

$$\langle U \rangle_{210} = \frac{eE_0}{32\pi a_0^5} \int_0^\infty r^3 e^{-r/a_0} dr \int_{-1}^{+1} \cos^3 \theta d \cos \theta \int_0^{2\pi} d\phi = 0 \quad (4)$$

$$\begin{aligned} \langle U \rangle_{21\pm 1} &= \frac{eE_0}{64(3\pi)a_0^5} \int_0^\infty r^3 e^{-r/a_0} dr \int_{-1}^{+1} \cos \theta \sin^2 \theta d \cos \theta \int_0^{2\pi} d\phi \\ &= \frac{eE_0}{96a_0} \int_0^\infty y^3 e^{-y} dy \int_{-1}^1 x(1-x^2) dx = 0, \quad y = \frac{r}{a_0}, \quad x = \cos \theta \quad (5) \end{aligned}$$

Atom in uniform magnetic field $\vec{B} = B_0 \hat{z}$, magnetic interaction with orbital magnetic moment of an electron is

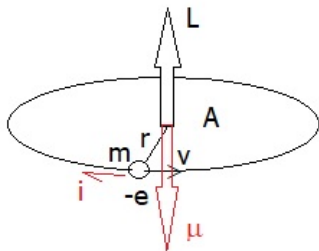
$$U_B = -\vec{\mu} \cdot \vec{B} = -\mu_z B_0 \quad (6)$$

where $\vec{\mu} = i\vec{A} = -\frac{ev}{2\pi r} \pi r^2 \hat{n} = -\frac{e}{2m_e} (m_e v r) \hat{n} = -\frac{e}{2m_e} \vec{L}$. Then we have from above

$$\vec{\mu} = \frac{-e}{2m_e} \vec{L}, \quad (7)$$

$$U_B = \frac{eB_0}{2m_e} L_z \quad (8)$$

$$\begin{aligned} \langle U_B \rangle_{nlm} &= \frac{eB_0}{2m_e} m\hbar \\ &= m\mu_B B_0 \end{aligned} \quad (9)$$

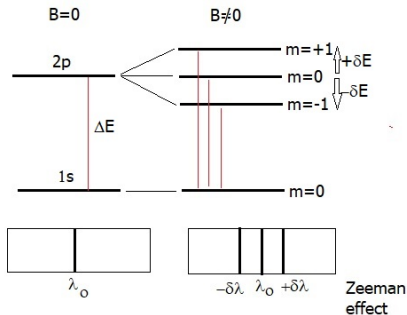


where $\mu_B = \frac{e\hbar}{2m_e} = 5.788 \times 10^{-5} \text{ eV/T}$ is Bohr magneton

Atomic transition with additional selection rule

$$\Delta m = 0, \pm 1$$

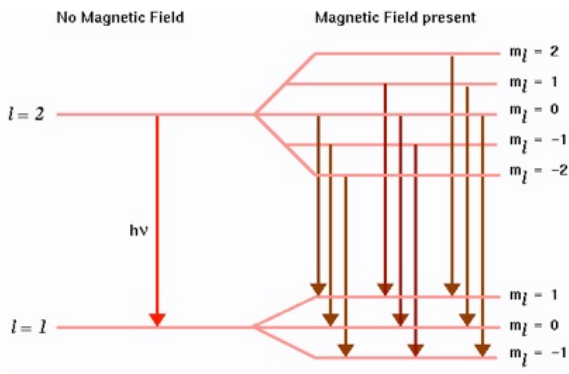
Spectral line splitting (Zeeman effect)



$$\lambda_0 = \frac{hc}{\Delta E} \mapsto \frac{\delta\lambda}{\lambda_0} = \frac{\delta E}{\Delta E} = \frac{\mu_B B_0}{\Delta E}$$

For example of L_α , $\Delta E = 10.2\text{eV}$, $\lambda_0 = 122.8\text{nm}$. With $B_0 = 1.0\text{T}$ we will have $\delta\lambda = 0.697 \times 10^{-3}\text{nm}$, the fourth decimal correction!

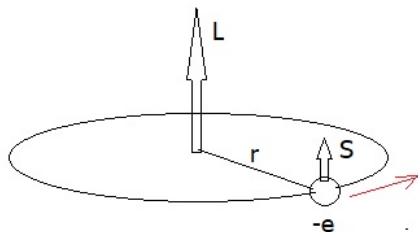
Other Zeeman splitting, i.e., H_{α} splitting



Can you estimate the spectral line splitting?

Fine structure

Spin-orbit coupling



Interaction potential energy is

$$\vec{\mu}_L = \frac{-e}{2m_e} \vec{L}, \quad \vec{\mu}_S = g \frac{-e}{2m_e} \vec{S}, \quad g = 2 \quad (10)$$

$$U_{LS} = -\vec{\mu}_S \cdot \vec{B}_L = g \frac{1}{2m_e^2 c^2} \frac{Ke^2}{r^3} (\vec{L} \cdot \vec{S}) \quad (11)$$

Effects on atomic spectrum

$$\langle U_{LS} \rangle_{nlm,sm_s} = g \frac{Ke^2}{2m_e^2 c^2} \langle \frac{1}{r^3} \rangle_{nl} \langle \vec{L} \cdot \vec{S} \rangle_{lm,sm_s} \quad (12)$$

Note that

$$\langle \frac{1}{r^3} \rangle_{nl} = \left(\frac{m_e c^2 \alpha}{\hbar c} \right)^3 \frac{1}{n^3} \frac{1}{l(l+1/2)l+1}, \quad l > 0 \quad (13)$$

$$\begin{aligned} \vec{J} &= \vec{L} + \vec{S} \mapsto J^2 = L^2 + S^2 + 2\vec{L} \cdot \vec{S} \\ \vec{L} \cdot \vec{S} &= \frac{1}{2} (J^2 - L^2 - S^2) \end{aligned} \quad (14)$$

with $s = 1/2, l = 0, 1, \dots, n-1$ and

$$J^2 = j(j+1)\hbar^2, j = l + \frac{1}{2}, l - \frac{1}{2}$$

Combine all together, we will have

$$\langle U_{LS} \rangle_{nlm,sm_s} = \frac{1}{4} m_e c^2 \alpha^4 \frac{1}{n^3} \frac{1}{j+1/2} \begin{cases} \frac{1}{j}, & j = l + 1/2 \\ -\frac{1}{j+1}, & j = l - 1/2 \end{cases} \quad (15)$$

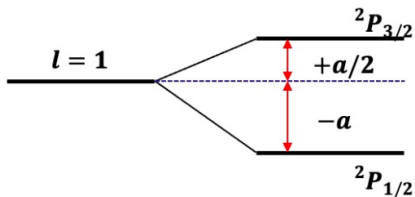


Figure 1.53: Splitting of $2p$ -state due to spin-orbit coupling.

Modified spectroscopic notation

$$nl \mapsto {}^nL_j$$

For example $1s \mapsto {}^1S_{1/2}$, $2p \mapsto {}^2P_{3/2}, {}^2P_{1/2}$ or $3d \mapsto {}^3D_{5/2}, {}^3D_{3/2}$

Exercise Estimate the LS-splitting of the d-orbital.