

# Lecture 5 Alpha Emission

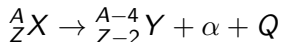
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# Alpha emission equation

- ▶  $\alpha = He^{+2}$  and  $M(\alpha) = 4.001506u = 3.7274MeV/c^2$
- ▶ Alpha emission equation



- ▶ Decay energy ( $K = p^2/2m$ )

$$Q = (M_X - M_Y - M_\alpha)c^2 = K_Y + K_\alpha \quad (1)$$

$$X - \text{rest frame : } p_\alpha = p_Y \mapsto M_\alpha K_\alpha = M_Y K_Y \quad (2)$$

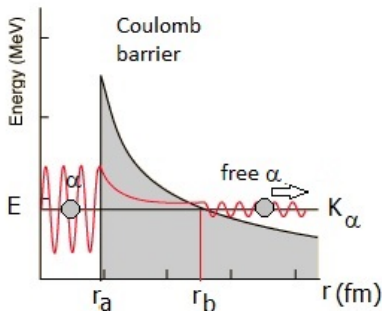
$$\frac{M_\alpha}{M_Y} = \frac{4}{A-4} \mapsto K_\alpha = \frac{A-4}{A}Q \quad (3)$$

For example  ${}^{214}_{84}Po \rightarrow {}^{210}_{82}Pb + \alpha$ , with  $Q = 7.83MeV$

$$K_\alpha = \frac{210}{214} \times 7.83MeV = 7.68MeV$$

# Gamow theory of alpha emission

- ▶ George Gamow (1928)
- ▶ Assumption: the  $\alpha$  particle is already exist inside the nucleus. Its emission is quantum tunneling through Coulomb barrier



- ▶ Tunneling coefficient, with  $Z = Z_Y$ ,

$$T(E) \simeq e^{-2 \int_{r_a}^{r_b} k(r) dr}, \quad k(r) = \sqrt{\frac{2M_\alpha}{\hbar^2} \left( \frac{2ZKe^2}{r} - E \right)}$$

- Calculation, using the fact that  $E = K_\alpha = \frac{2ZK_e^2}{r_b}$ , then we can write

$$\ln T = -2\sqrt{\frac{2M_\alpha E}{\hbar^2}} \int_{r_a}^{r_b} \sqrt{\frac{r_b}{r} - 1} dr \quad (4)$$

$$\text{Assign : } x = r_b \cos^2 \theta, \quad \theta = \cos^{-1} \sqrt{\frac{x}{r_b}}$$

$$dx = -2r_b \sin \theta \cos \theta d\theta$$

$$\mapsto \int \sqrt{\frac{r_b}{r} - 1} dr = -2r_b \int \sin^2 \theta d\theta = -2r_b \left[ \frac{\theta}{2} - \frac{1}{4} \sin(2\theta) \right]$$

$$\mapsto \ln T = -2\sqrt{\frac{2M_\alpha E}{\hbar^2}} r_b \left[ \cos^{-1} \sqrt{\frac{r_a}{r_b}} - \sqrt{\frac{r_a}{r_b}} \sqrt{1 - \frac{r_a}{r_b}} \right] \quad (5)$$

In the limit of wide barrier,  $r_b \gg r_a$ , we will have

$$\cos(\pi/2 - \theta) = \sin(\theta) \xrightarrow{\theta \rightarrow 0} \theta$$

$$\cos \left( \frac{\pi}{2} - \sqrt{\frac{r_a}{r_b}} \right) \simeq \sqrt{\frac{r_a}{r_b}} \mapsto \cos^{-1} \sqrt{\frac{r_a}{r_b}} \simeq \frac{\pi}{2} - \sqrt{\frac{r_a}{r_b}} \quad (6)$$

- Calculation (cont.), with  $1 - r_a/r_b \simeq 1$ , we have

$$\ln T \simeq -2\sqrt{\frac{2M_\alpha E}{\hbar^2}} r_b \left[ \frac{\pi}{2} - \sqrt{\frac{r_a}{r_b}} \right] \quad (7)$$

Replacing  $r_b = \frac{2ZKe^2}{E}$ , we will have

$$\ln T \simeq 8\sqrt{\frac{M_\alpha Ke^2}{\hbar^2}} Z^{1/2} r_a^{1/2} - 4\pi Ke^2 \sqrt{\frac{M_\alpha}{2\hbar^2}} ZE^{-1/2} \quad (8)$$

$$\mapsto 2.97 Z^{1/2} r_a^{1/2} - 3.95 ZE^{-1/2} \quad (9)$$

with  $[E] = \text{MeV}$ ,  $[r_a] = \text{fm}$

- Using the fact that  $\ln A = \frac{1}{\log_{10} e} \log_{10} A = 0.4343 \log_{10} A$ , we will have from (9)

$$\log_{10} T \simeq 1.29 Z^{1/2} r_a^{1/2} - 1.72 ZE^{-1/2} \quad (10)$$

## Decay constant

- ▶ Frequency of the alpha particle comes to hit the Coulomb barrier is determined to be

$$f = \frac{v_\alpha}{2r_a}, \quad v_\alpha = \sqrt{\frac{2K_\alpha}{M_\alpha}}$$

- ▶ The decay constant becomes

$$\lambda = fT = \frac{v_\alpha}{2r_a} T \quad (11)$$

$$\log_{10} \lambda \simeq \log_{10} \frac{v_\alpha}{2r_a} + 1.29Z^{1/2}r_a^{1/2} - 1.72ZE^{-1/2} \quad (12)$$

This is known as *Geiger-Nuttall law*:

$$\log_{10} \lambda \simeq -a_1 \frac{Z}{\sqrt{E}} + a_2, \quad \log_{10} T_{1/2} \simeq a_1 \frac{Z}{\sqrt{E}} - a'_2$$

► Geiger-Nuttall plot

