

Lecture 5 Alpha Emission

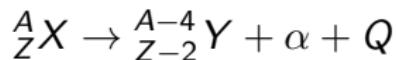
SCPY327 ANP, Physics-MUSC, 2021/2

Udom Robkob

March 24, 2022

Alpha emission equation

- ▶ $\alpha = He^{+2}$ and $M(\alpha) = 4.001506u = 3.7274MeV/c^2$
- ▶ Alpha emission equation



- ▶ Decay energy ($K = p^2/2m$)

$$Q = (M_X - M_Y - M_\alpha)c^2 = K_Y + K_\alpha \quad (1)$$

$$X - \text{rest frame} : \quad p_\alpha = p_Y \mapsto M_\alpha K_\alpha = M_Y K_Y \quad (2)$$

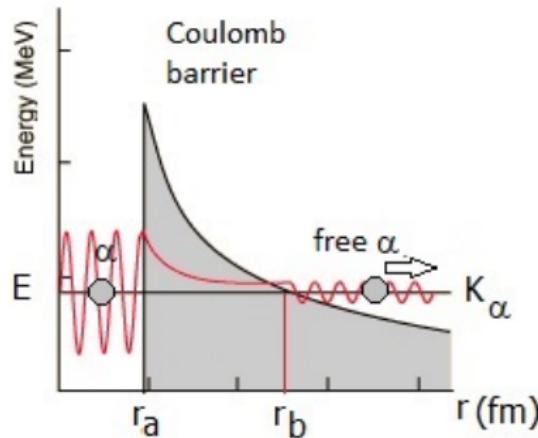
$$\frac{M_\alpha}{M_Y} = \frac{4}{A-4} \mapsto K_\alpha = \frac{A-4}{A} Q \quad (3)$$

For example $^{214}_{84}Po \rightarrow ^{210}_{82}Pb + \alpha$, with $Q = 7.83MeV$

$$K_\alpha = \frac{210}{214} \times 7.83MeV = 7.68MeV$$

Gamow theory of alpha emission

- ▶ George Gamow (1928)
- ▶ Assumption: the α particle is already exist inside the nucleus. Its emission is quantum tunneling through Coulomb barrier



- ▶ Tunneling coefficient, with $Z = Z_Y$,

$$T(E) \simeq e^{-2 \int_{r_a}^{r_b} k(r) dr}, \quad k(r) = \sqrt{\frac{2M_\alpha}{\hbar^2} \left(\frac{2ZKe^2}{r} - E \right)}$$

- ▶ Calculation, using the fact that $E = K_\alpha = \frac{2ZKe^2}{r_b}$, then we can write

$$\ln T = -2 \sqrt{\frac{2M_\alpha E}{\hbar^2}} \int_{r_a}^{r_b} \sqrt{\frac{r_b}{r} - 1} dr \quad (4)$$

Assign : $x = r_b \cos^2 \theta, \theta = \cos^{-1} \sqrt{\frac{x}{r_b}}$
 $dx = -2r_b \sin \theta \cos \theta d\theta$

$$\begin{aligned} \mapsto \int \sqrt{\frac{r_b}{r} - 1} dr &= -2r_b \int \sin^2 \theta d\theta = -2r_b \left[\frac{\theta}{2} - \frac{1}{4} \sin(2\theta) \right] \\ \mapsto \ln T &= -2 \sqrt{\frac{2M_\alpha E}{\hbar^2}} r_b \left[\cos^{-1} \sqrt{\frac{r_a}{r_b}} - \sqrt{\frac{r_a}{r_b}} \sqrt{1 - \frac{r_a}{r_b}} \right] \end{aligned} \quad (5)$$

In the limit of wide barrier, $r_b \gg r_a$, we will have

$$\cos(\pi/2 - \theta) = \sin(\theta) \xrightarrow{\theta \rightarrow 0} \theta$$

$$\cos \left(\frac{\pi}{2} - \sqrt{\frac{r_a}{r_b}} \right) \simeq \sqrt{\frac{r_a}{r_b}} \mapsto \cos^{-1} \sqrt{\frac{r_a}{r_b}} \simeq \frac{\pi}{2} - \sqrt{\frac{r_a}{r_b}} \quad (6)$$

- ▶ Calculation (cont.), with $1 - r_a/r_b \simeq 1$, we have

$$\ln T \simeq -2 \sqrt{\frac{2M_\alpha E}{\hbar^2}} r_b \left[\frac{\pi}{2} - \sqrt{\frac{r_a}{r_b}} \right] \quad (7)$$

Replacing $r_b = \frac{2ZKe^2}{E}$, we will have

$$\ln T \simeq 8 \sqrt{\frac{M_\alpha Ke^2}{\hbar^2}} Z^{1/2} r_a^{1/2} - 4\pi Ke^2 \sqrt{\frac{M_\alpha}{2\hbar^2}} Z E^{-1/2} \quad (8)$$

$$\mapsto 2.97 Z^{1/2} r_a^{1/2} - 3.95 Z E^{-1/2} \quad (9)$$

with $[E] = MeV$, $[r_a] = fm$

- ▶ Using the fact that $\ln A = \frac{1}{\log_{10} e} \log_{10} A = 0.4343 \log_{10} A$, we will have from (9)

$$\log_{10} T \simeq 1.29 Z^{1/2} r_a^{1/2} - 1.72 Z E^{-1/2} \quad (10)$$

Decay constant

- ▶ Frequency of the alpha particle comes to hit the Coulomb barrier is determined to be

$$f = \frac{v_\alpha}{2r_a}, \quad v_\alpha = \sqrt{\frac{2K_\alpha}{M_\alpha}}$$

- ▶ The decay constant becomes

$$\lambda = fT = \frac{v_\alpha}{2r_a} T \quad (11)$$

$$\log_{10} \lambda \simeq \log_{10} \frac{v_\alpha}{2r_a} + 1.29 Z^{1/2} r_a^{1/2} - 1.72 Z E^{-1/2} \quad (12)$$

This is known as *Geiger–Nuttall law*:

$$\log_{10} \lambda \simeq -a_1 \frac{Z}{\sqrt{E}} + a_2, \quad \log_{10} T_{1/2} \simeq a_1 \frac{Z}{\sqrt{E}} - a'_2$$

► Geiger-Nuttall plot

