Lecture 6 Beta Emission SCPY327 ANP, Physics-MUSC, 2021/2

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Today topics

- ► Fermi golden rule
- ▶ Beta emission

Fermi golden rule

Basic Schrodinger equation

$$\hat{H}_0|n\rangle = E_n|n\rangle$$

With time-dependent interaction $H(t) = H_0 + V(t)$, the Schrodinger equation in interaction picture is

$$i\hbar\partial_t |\psi(t)\rangle_S = \hat{H}(t)|\psi(t)\rangle_S, \ |\psi(t)\rangle_S = \sum_n c_n(0)e^{-iE_nt/\hbar}|n\rangle \ (1)$$

$$|\psi(t)\rangle_I = e^{iH_0t/\hbar}|\psi(t)\rangle_S, \ V_I(t) = e^{iH_0t/\hbar}V(t)e^{-iH_0t/\hbar}$$
 (2)

$$\mapsto i\hbar\partial_t |\psi(t)\rangle_I = V_I(t)|\psi(t)\rangle_I \quad (3)$$

$$|\psi(t)\rangle_I = U_I(t,0)|\psi(0)\rangle_I \mapsto U_I(t,0) = e^{-i\int_0^t V_I(t')dt'/\hbar}$$
 (4)

$$|\psi(t)\rangle_I = \sum_n |n\rangle \underbrace{\langle n|U_I(t,0)|\psi(0)\rangle_I}_{c_n(t)}$$
 (5)

► Perturbation theory

$$c_n^{(1)}(t) = -\frac{i}{\hbar} \int_0^t \langle n|V_I(t')|m\rangle dt'$$
$$= -\frac{i}{\hbar} \int_0^t V_{nm}(t')e^{i\omega_{nm}t'}dt'$$
(8)

$$P_n(t) = |c_n(t)|^2 \simeq |c_n^{(1)}(t)|^2$$
 (9)

► Harmonic perturbation

$$V(t) = Ve^{-i\omega t} \mapsto c_f^{(1)}(t) = -\frac{i}{\hbar} \int_0^t V_{fi} e^{i(\omega_{fi}t - \omega)t'} dt'$$
$$= -\frac{i}{\hbar} V_{fi} \frac{e^{i(\omega_{fi} - \omega)t} - 1}{i(\omega_{fi} - \omega)} \quad (10)$$

$$P_{i o f} = rac{1}{\hbar^2} |V_{fi}|^2 \left(rac{\sin((\omega_{fi} - \omega)t/2)}{(\omega_{fi} - \omega)/2}
ight)^2 \simeq rac{2\pi t}{\hbar^2} |V_{fi}|^2 \delta(\omega_{fi} - \omega)$$

▶ The first Fermi golden rule of transition rate

$$W_{i\to f} = \frac{P_{i\to f}}{t} = \frac{2\pi}{\hbar^2} |v_{fi}|^2 \delta(\omega_{fi} - \omega)$$
 (11)

lacktriangle Transition to a group of final states with density of state $ho(E_f)$

$$W_{i\to f} = \frac{2\pi}{\hbar^2} |V_{fi}|^2 \int \rho(E_F) \delta(\omega_{fi} - \omega) dE_F \qquad (12)$$

$$\simeq \frac{2\pi}{\hbar} |Vfi|^2 \rho(E_F) \tag{13}$$

The second Fermi golden rule.

Beta emission

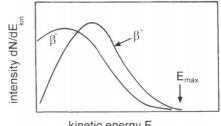
- $\beta^{\pm} = \pm e$
- Decay equations

$$_{Z}^{A}X \to _{Z+1}^{A}Y + \beta^{-} + \bar{\nu}_{e} + Q_{\beta^{-}}$$

 $_{Z}^{A}X \to _{Z-1}^{A}Y + \beta^{+} + \nu_{e} + Q_{\beta^{+}}$

These come from fundamental decays $n^0 o p^+1 + \beta^- + \bar{\nu}_e$ and $p^+ o n^0 + \beta^+ + \nu_e$

▶ Beta spectrum and the existence of neutrino/anti-neutrino (W.Pauli, 1930)



kinetic energy E_{kin}



- ► Fermi theory of beta emission
- The number of state of fermion having momentum between $p \rightarrow p + dp$ inside a volume V is

$$dN = \frac{V}{2\pi^2\hbar^3} p^2 dp \mapsto \frac{dN}{dE} = \frac{V}{2\pi^2\hbar^3} p^2 \frac{dp}{dE}$$

For the neutrino

$$p_{\nu} = E_{\nu}/c \mapsto dp_{\nu} = \frac{dE_{\nu}}{c} \mapsto \frac{dN_{\nu}}{dE_{\nu}} = \frac{V}{2\pi^2\hbar^3c^3}E_{\nu}^2$$
 (14)

For the beta

$$p_e c = \sqrt{E_e^2 - (m_e c^2)^2} \mapsto dp_e c = \frac{E_e dE_e}{\sqrt{E_e^2 - (m_e c^2)^2}}$$
 (15)

$$\mapsto \frac{dN_e}{dE_e} = \frac{V}{2\pi^2\hbar^3c^3}\sqrt{E_e^2 - (m_ec^2)^2}E_e \quad (16)$$

$$\mapsto \frac{dN_{\rm e}}{dK_{\rm e}} = \frac{V}{2\pi^2\hbar^3c^3} \sqrt{K_{\rm e}(K_{\rm e} + 2m_{\rm e}c^2)} (K_{\rm e} + m_{\rm e}c^2) \tag{17}$$

with the fact that $E_e = K_e + m_e c^2$



From Fermi golden rule, with neutrino momentum q

$$\begin{split} \frac{dW_{e,fi}}{dp_{e}} &= \frac{2\pi}{\hbar} |V_{fi}|^{2} \frac{dN_{e}}{dp_{e}} \frac{dN_{\nu}}{dE_{fi}} (18) \\ E_{fi} &= K_{e} + cq \equiv Q \mapsto dE_{fi} = cdq \; (\textit{fixed} \; K_{e}) (19) \\ \mapsto \frac{dW_{e,fi}}{dp_{e}} &= \frac{2\pi}{\hbar} |V_{fi}|^{2} \frac{V}{2\pi^{2}\hbar^{3}} p_{e}^{2} \frac{V}{2\pi^{2}\hbar^{3}} q^{2} \delta(E_{fi} - (K_{e} + K_{\nu})) (20) \end{split}$$

The delta function is inserted to account the energy conservation.

Let $M_{fi} = VV_{fi}$ and $qc = (Q - K_e)$, after integrate (20) over all neutrino momenta using delta function, we will have

$$\frac{dW_{e,fi}}{dp_{e}} = \frac{|M_{fi}|^{2}}{2\pi^{3}\hbar^{7}c^{3}}p_{e}^{2}(Q - K_{e})^{2}$$

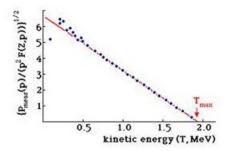
This expression was first derived by Fermi, and display the beta spectrum.

More details analysis of M_{fi} from weak theory will lead Fermi expression to

$$rac{dW_{e,fi}}{dp_e} \propto |M_{fi}|^2 F(Z,p_e) p_e^2 (Q-K_e)^2$$
 $N_e \simeq F(Z,p_e) p_e^2 (Q-K_e)^2$

where $F(Z, p_e)$ is known as Fermi function.

Kurie plot of beta spectrum



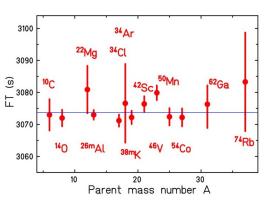
Decay constant

$$\lambda = \int dW_{e,fi} = \frac{1}{2\pi^3 \hbar^7 c^3} \int_0^{\rho_{e,max}} |M_{fi}|^2 F(Z,Q) p_e^2 (Q - K_e)^2 dp_e$$

$$\mapsto \lambda \simeq \frac{|M_{fi}|^2}{2\pi^3 \hbar^7 c^3} f(Z,Q) \quad \text{"Fermi integral" (21)}$$

$$\mapsto ft_{1/2} = 0.693 \frac{2\pi^3 \hbar^7 c^3}{|M_{fi}|^2}$$
 "comparative half life" (22)

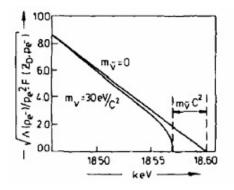
► ft-value





Neutrino mass

▶ The existence of neutrino mass appear in Kurie spectrum



$$N_{\rm e} \propto F(Z, p_{\rm e}) p_{\rm e}^2 (Q - K_{\rm e}) \sqrt{(Q - K_{\rm e})^2 - m_{\nu}^2 c^4}$$

Parity violation

- ➤ Tsung-Dao Lee and Chen-Ning Yang, the theoretical physicists who originated the idea of parity nonconservation and proposed the experiment, received the 1957 Nobel Prize in physics for this result. Chien-Shiung Wu's role in the discovery was mentioned in the Nobel prize acceptance speech, but was not honored until 1978, when she was awarded the first Wolf Prize.
- Parity violation

