Lecture 11 Nuclear Radiations ICPY473 Nuclear Physics, MUIC-3-Trimester, 2020-21

Udom Robkob, Physics-MUSC

Monday 7, JUne 2021

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Lecture 11 Nuclear Radiations

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Today Topics

- The radiations
- Interaction of photons with matter
- Interaction of charged particles with matter
- Radiations detection and measurement
- Radiations protection

The Radiations

- Nuclear radiations are energetic particles and gamma photon from nuclear processes
- They are classified into
 - photon, i.e., X-rays, γ-ray
 - charged particles, i.e, β^{\pm} , p^{\pm} , α , heavy ions
 - charge neutral particles, i.e., neutron and neutrino
- $\bullet\,$ Three of them, α,β,γ come from decay of radioisotopes, and are common to us
- We will study their interactions with matter, which will lead us to their detection, measurement and protection

Interaction of Photon with Matter

- Photons are quantum of electromagnetic waves, they interaction with mater with electromagnetic force, this results to *ionization* of matter atoms
- The main mechanisms for energy lose are
 - photoelectric effects
 - Compton scattering
 - pair production

Interaction of Photon with Matter

- Photons are quantum of electromagnetic waves, they interaction with mater with electromagnetic force, this results to *ionization* of matter atoms
- The main mechanisms for energy lose are
 - photoelectric effects
 - Compton scattering
 - pair production
- The *cross section* of low energy photon photoelectric interaction with matter atom is

$$\sigma_{ph} = \begin{cases} Z^4 / (h\nu)^3, & \text{low energy} \\ Z^5 / h\nu, & \text{high energy} \end{cases}$$
(1)

• Photoelectric cross section



Figure: H. Hirayama, KEK-report, 2000

• The *cross section* of moderate energy photon Compton scattering of atomic electron of matter atom, according to *Klein-Nishina*, is

$$\sigma_{KN} = 2\pi r_0^2 \left\{ \frac{1+k}{k^2} \left[\frac{2(1+k)}{1+2k} - \frac{\ln(1+2k)}{k} \right] + \frac{\ln(1+2k)}{2k} - \frac{1+3k}{(1+2k)^2} \right\}$$
(2)

where $r_0 = \frac{e^2}{4\pi\epsilon_0^2 m_e c^2} \simeq 2.818 fm$ is classical electron radius, and $k = \frac{h\nu}{m_e c^2}$ is dimensionless photon kinematic variable.

• Thomson limit, $k \rightarrow 0$, of Klein-Nishina formula is

$$\sigma_T = \frac{8\pi}{3} r_0^2 \tag{3}$$

• Compton scattering cross section of photon



Figure: O. Klein and Y. Nishina, Z. Physik (52), 1929 (853-868).

• The *cross section* of energetic photon pair production, according to Bethe-Heitler formula, is

$$\sigma_{pair}(E_{\gamma}) = \simeq \alpha Z^2 r_0^2 \ln E_{\gamma} \tag{4}$$

where α is fine structure constant and Z is nuclear charge of target atom.

According to the pair production process

$$\gamma + X \rightarrow X + 2e^- + Q$$

It appears with threshold energy of pair production as

$$T_{th}^{CM} = 2m_e c^2 \tag{5}$$

$$Q = -2m_e c^2 (-T_{th}^{CM})$$
 (6)

$$T_{th}^{LAB} = \frac{2m_e c^2 (m_e c^2 + M_X c^2)}{M_X c^2}$$
(7)

• Pair production cross section of photon



Figure: H. Bethe and W. Heitler (1954).

• Absorption of photon by matter and linear attenuation coefficient μ

$$dI(x) \propto -dx \rightarrow dI(x) = -\mu dx \rightarrow I(x) = I(0)e^{-\mu x}$$
(8)

$$\mu = n\sigma \quad [m^{-1}] \tag{9}$$

where $n[\#/m^3]$ is number density of target atom and $\sigma[m^2]$ is the cross section of the corresponding interaction.



• Mass attenuation coefficient μ_m is defined with mass density $\rho[{\rm gm}/{\rm cm}^3]$ as

$$\mu_m = \left(\frac{\mu}{\rho}\right)\rho \to [cm^2/gm] \tag{10}$$

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• Mass attenuation coefficient of water



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• Table of mass attenuation coefficient of some matters

TABLE 24.1

Mass attenuation coefficient as a function of X-ray photon energy in water, air, muscle, bone, and fat. Average densities, atomic numbers and electron number densities are included, with values for air at p = 1.0 atm.

| | Water | Air | Muscle | Fat | Bone |
|-------------------------------------|----------------------------|-------------------------|----------------------------|----------------------------|----------------------------|
| | μ/ρ | μ/ρ | μ/ρ | μ/ρ | μ/ρ |
| Energy (keV) | (cm ² /g) | (cm ² /g) | (cm²/g) | (cm²/g) | (cm ² /g) |
| 5 | 42.1 | 40.3 | 42.0 | 24.6 | 138 |
| 10 | 5.07 | 4.91 | 5.15 | 3.08 | 19.8 |
| 15 | 1.57 | 1.52 | 1.60 | 1.01 | 6.19 |
| 20 | 0.761 | 0.733 | 0.778 | 0.533 | 2.75 |
| 30 | 0.361 | 0.340 | 0.365 | 0.296 | 0.953 |
| 40 | 0.263 | 0.243 | 0.264 | 0.235 | 0.509 |
| 50 | 0.224 | 0.205 | 0.224 | 0.210 | 0.347 |
| 80 | 0.183 | 0.166 | 0.182 | 0.179 | 0.208 |
| 100 | 0.171 | 0.154 | 0.169 | 0.168 | 0.180 |
| 150 | 0.151 | 0.136 | 0.149 | 0.150 | 0.149 |
| 200 | 0.137 | 0.123 | 0.136 | 0.137 | 0.133 |
| 1000 | 0.071 | 0.064 | 0.070 | 0.071 | 0.068 |
| 5000 | 0.030 | 0.028 | 0.030 | 0.030 | 0.030 |
| Density (kg/m ³) | 1000 | 1.20 | 1040 | 915 | 1650 |
| Electron density (<i>e</i> /kg) | 3.34 × 10 ²⁶ | 3.01 × 10 ²⁶ | 3.31 × 10 ²⁶ | 3.34 × 10 ²⁶ | 3.19 × 10 ²⁶ |
| Atomic number <z></z> | 7.5 | 7.8 | 7.6 | 6.5 | 12.3 |

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• The halve value thickness (HVT) $d_{1/2}$ is defined as

$$I(d_{1/2}) = \frac{1}{2}I_0 = I_0 e^{-\mu d_{1/2}} \to d_{1/2} = \frac{\ln 2}{\mu} \quad [cm]$$
 (11)

• It seem that HVT is matter dependent.

| Absorber | 100 keV | 200 keV | 500 keV |
|-----------|---------|---------|---------|
| Air | 3555 | 4359 | 6189 |
| Water | 4.15 | 5.1 | 7.15 |
| Carbon | 2.07 | 2.53 | 3.54 |
| Aluminium | 1.59 | 2.14 | 3.05 |
| Iron | 0.26 | 0.64 | 1.06 |
| Copper | 0.18 | 0.53 | 0.95 |
| Lead | 0.012 | 0.068 | 0.42 |

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• The *halve value layer* (HVL) is then defined by mutiplical with mass density of an absorber

$$HVL = \rho d_{1/2}$$

in unit of $[gm/cm^2]$, and become matter independent and it is characteristic to the photon. It is an amount of matter in gram per square centimeter to reduce photon density by one-half.

Interaction of Charged Particle with Matter

- The interaction is electromagnetic interaction, and results with excitation or ionization of medium atom. They are three types of charged particle energy lost
 - elastic scattering
 - inelastic scattering
 - Bremsstrahlung



• Stopping power *S*(*E*) is defined to be the energy lose rate of charged particle inside matter as

$$S(E) = -\frac{dE}{dx} \tag{12}$$

• *Range R* is defined to be the longest of penetration depth of struggling path of the charged particle inside matter.



• By definition

$$R = \int_{E}^{0} (-dE/dx) dE = \int_{E}^{0} S(E) dE$$
(13)

Bethe-Bloch formula

$$-\frac{dE}{dx} = \frac{4\pi}{m_e c^2} \cdot \frac{n z^2 (k_e e^2)^2}{\beta^2} \cdot \left[\ln\left(\frac{2m_e c^2 \beta^2}{I(1-\beta^2)} - \beta^2\right) \right]$$
(14)

where

$$n = \frac{N_A \cdot Z \cdot \rho}{A \cdot M}$$

is the electron density of the medium, $\beta = v/c$ is a particle velocity (in unit of light speed) and charged z. I is the ionization energy.

• Bethe-Bloch stopping power



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• The range and stopping power are related by Bragg curve



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• Bragg curve at various energies



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Radiations Detection and Measurement

Radiation detectors



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• Gas-filled detectors





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• Solid state detectors



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• Solid state detectors



• Charge-coupling (capacitors) detector



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• Radiations measurement

| lonizing | Roentgen | Charge/unit | 1 R = 2.58 × 10 ⁻⁴ C/kg |
|--|------------------------|-------------------------------------|------------------------------------|
| radiation | (R) | mass | |
| Absorbed | Rad | 1 rad = 0.01 J/kg | 1 rad = 0.01 Gy |
| dose | Gray (Gy) | 1 Gy = 1 J/kg | 1 Gy = 100 rad |
| Equivalent dose/ effective dose | Rem Sievert (Sv) | $Rem = rad \times W$ Sv = Gy × W | 1 rem = 0.01 Sv 1 Sv = 100 rem |

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• Radiations measurement

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• Weigh factor

| Radiation type and energy range | Radiation weighting factor, wR |
|--|--------------------------------|
| Photons, all energies | 1 |
| Electrons and muons, all energies | 1 |
| Protons and charged pions | 2 |
| Alpha particles, fission fragments, heavy nuclei | 20 |

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Radiations Protections

• Follows the rule of three *L*:

Long distance, Large shielding, Least time



Radiations Protections

• Follows the rule of three *L*:

Long distance, Large shielding, Least time



Limit Time



Increase Distance



• Beware of radiations sign



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