Lecture 23 Aspects of New Physics after LHC ICPY473 Nuclear Physics, MUIC, 3-Trimester, 2021

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Today Topics

- Higgs physics
- Beyond the standard model physics

Higgs Physics

- Why do we need Higgs particle? According to the similarity of QED and Weak interaction, except the massless photon γ and the massive vector bosons W[±], Z⁰.
- Energy range of QED in MeV, while weak interaction is GeV, so that something seem to be continuity from weak to QED by lowering the energy scale
- In classical physics, lowering in energy results to phase transition, i.e., gases condense into liquids and crystallize into solids, and all phase transitions are described by symmetry breaking mechanism



- Goldstone's theorem: for breaking of any continuous symmetry there will correspond with massless boson, it is called Goldstone's boson
- For magnetic materials, there are randomly orientation of magnetic moments (higher symmetry) after magnetization (directional oriented of magnetic moments) there appear with magnon as its elementary excitation
- For Helium gas, after Bose condensation there appear with rotons as its elementary excitation



Higgs mechanism: for breaking of any continuous symmetry there will correspond with massive boson, it is called Higgs's boson



With $D_{\mu} = \partial_{\mu} - iqA_{\mu}$ and $F_{\mu\nu} = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu}$, Higgs Lagrangian is

$$\mathcal{L} = [D_{\mu}\phi]^*[D^{\mu}\phi] - \mathcal{V}(\phi) - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$$

where $\mathcal{V} = -\mu^2 \phi^* \phi + \lambda (\phi^* \phi)^2$ is symmetry breaking potential

Gauge boson get mass from Higgs as

$$-\frac{1}{4}F_{\mu\nu}F^{\mu\nu} \to -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \frac{1}{2}m^2A_{\mu}A^{\mu}, \quad m^2 = g^2v^2$$

where $v=\langle \phi \rangle$ is the vacuum expectation value of the Higgs field

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Where is Higgs? Higgs production phenomenology



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The Higgs production cross section



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> The Higgs existence is observed through its decay remnants



Higgs detection at CMS-CERN



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Higgs detection at ATLAS-CERN



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Celebration on October 2, 2015



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Beyond the Standard Model Physics

Problems with standard model

Too many parameters

me	Electron	mass	511	keV		
mμ	Muon mas	s	105.7 M	eV		
mτ	Tau mass	5	1.78 Ge	v		
mu	Up quark	mass	μ MS = 2	GeV	1.9 MeV	
md	Down qua	ark mass	μMS	= 2 Ge	v 4.4	MeV
ms	Strange	quark m	ass µMS	= 2 Ge	V 87	MeV
mc	Charm qu	uark mas	s µMS	= mc	1.32 Ge	v
mb	Bottom d	quark ma	ss µMS	= mb	4.24 Ge	v
mt	Top quar	rk mass	On-shel	1 schem	e 172	.7 GeV
012	CKM	12-mixi	ng angle		13.1°	
823	СКМ	23-mixi	ng angle		2.4°	
013	CKM	13-mixi	ng angle		0.2°	
δ	CKM CP-	violatin	g Phase	0.	995	
g1	or g'	U(1) ga	uge coup	ling	µMS = m	z 0.35
g2	or g	SU(2) g	auge cou	pling	µMS = m	Z 0.65
g3	or gs	SU(3) g	auge cou	pling	µMS = m	z 1.22
BQC	D QCD	vacuum	angle	~0		
v	Higgs va	acuum ex	pectatio	n value	246	GeV
mH	Higgs ma	ass	~ 125 G	eV (ten	tative)	

• $\alpha_1 = (5/3)g'^2/(4\pi) = 5\alpha/(3\cos^2\theta_W)$ • $\alpha_2 = g^2/(4\pi) = \alpha/\sin^2\theta_W$ • $\alpha_3 = g_s^2/(4\pi)$ • $\alpha_1(M_Z) = 0.017$ • $\alpha_2(M_Z) = 0.034$

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•
$$\alpha_3(M_Z) = 0.118 \pm 0.003$$

• Mass hierarchy problem of δm^2 .

Problems with standard model (cont.)

Asymmetry of particles and anti-particles



Figure 2. Asymmetric decay of a symmetric state.

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Problems with standard model (cont.)

Asymmetry of particles and anti-particles



Beyond the standard model





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supersymmetry unification



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Extra dimensions, braneworld extension



Randall-Sundrum model



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