

4 Many-Nucleon Models

4.1 Fermi gas model

For $N + Z$ nucleon system, from Fermi statistics, the system energy is

$$\epsilon_k = \frac{\hbar^2 k^2}{2M} \rightarrow E = \sum_{k \in [0, k_F]} \epsilon_k \quad (1)$$

From basic quantum mechanics of infinite 3D square well, we have

$$\vec{k} = (k_x = n_x \pi / L, n_y \pi / L, n_z \pi / L) \quad (2)$$

$$\epsilon_k \rightarrow \epsilon_{n_x n_y n_z} = \frac{\hbar^2}{2ML^2} (n_x^2 + n_y^2 + n_z^2) = \frac{\hbar^2 n^2}{2ML^2} \quad (3)$$

$$\vec{n} = (n_x, n_y, n_z) \rightarrow dn = \frac{1}{8} \frac{L^3}{\pi^3} d\vec{k} = \frac{V}{(2\pi\hbar)^3} d\vec{p}, \quad \vec{p} = \hbar\vec{k} \quad (4)$$

$$n = \frac{V}{(2\pi\hbar)^3} (4\pi) \int_0^{p_F} p^2 dp = \frac{V}{(2\pi\hbar)^3} \frac{4\pi p_F^3}{3} = \frac{V p_F^3}{6\pi^2 \hbar^3} \quad (5)$$

$$\text{With spin : } \rightarrow n = (2) \frac{V p_F^3}{6\pi^2 \hbar^3} = \frac{V p_F^3}{3\pi^2 \hbar^3} \quad (6)$$

Then we have

$$N = \frac{V (p_F^n)^3}{3\pi^2 \hbar^3}, \quad Z = \frac{V (p_F^p)^3}{3\pi^2 \hbar^3} \quad (7)$$

Let us determine

$$r = r_0 A^{1/3} \rightarrow V = \frac{4\pi r^3}{3} = \frac{4\pi r_0^3 A}{3} \quad (8)$$

$$\rightarrow n = \frac{4\pi r_0^3 A}{3} \cdot \frac{p_F^3}{3\pi^2 \hbar^3} = \frac{4A}{9\pi} \frac{r_0^2 p_F^3}{\hbar^3} \rightarrow p_F = \left(\frac{9\pi n}{4A} \right)^{1/3} \frac{\hbar}{r_0} \quad (9)$$

In a simple case of $n = N = Z = A/2$, we will have

$$p_F = \left(\frac{9\pi}{8} \right)^{1/3} \frac{\hbar}{r_0} \simeq 250 \text{ MeV}/c \rightarrow E_f = \frac{p_F^2}{2M} \simeq 33 \text{ MeV} \quad (10)$$

The depth of nuclear potential well $-U_0$ can be estimated to be

$$U_0 = E_F + e_b \simeq 40 \text{ MeV}$$

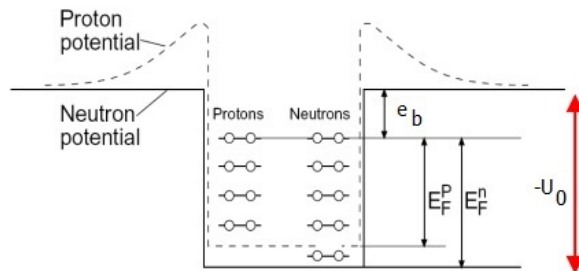


Figure 1:

4.2 Liquid drop model

In order to explain nuclear fission, the liquid drop nuclear model was proposed by George Gamow (1930)

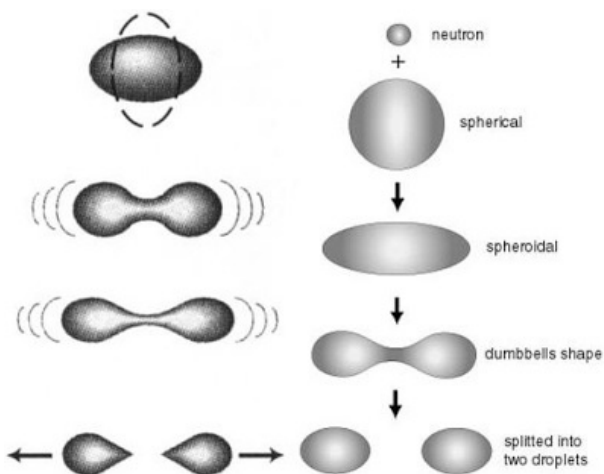


Figure 2:

The nuclear binding energy formula, determined from the liquid drop, was proposed by Carl Friedrich von Weizsäcker (1935), and known as *Bethe-Weizsäcker formula*, in the form

$$E_b = a_V \cdot A - a_S \cdot A^{2/3} - a_C \cdot \frac{Z^2}{A^{1/3}} - a_A \cdot \frac{(Z - N)^2}{A} - \delta \cdot \frac{1}{A^{1/2}} \quad (11)$$

$$a_V \simeq 16\text{MeV}, a_S \simeq 20\text{MeV}, a_C \simeq 0.75\text{MeV}, a_A \simeq 21\text{MeV} \quad (12)$$

$$\delta = \begin{cases} -11.2\text{MeV}, & \text{even} - Z, \text{even} - N \\ 0, & \text{odd} - A \\ +11.2\text{MeV}, & \text{odd} - Z, \text{odd} - N \end{cases} \quad (13)$$

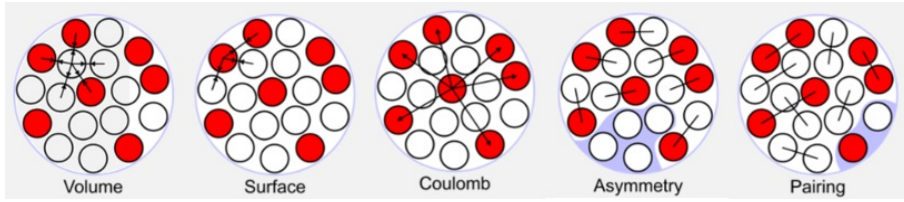


Figure 3:

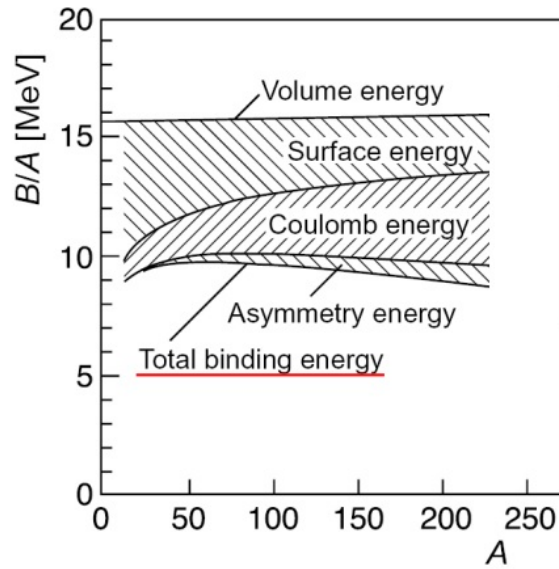


Figure 4:

4.3 Deformed sphere (Nilsson) model

Nuclear shape can be described by the surface function

$$R(\theta, \phi) = R_0 \left(1 + \sum_l \sum_{m=-l}^{+l} \alpha_{lm} Y_{lm}^*(\theta, \phi) \right) \quad (14)$$

$$l = 0 \text{ compression} \quad (15)$$

$$l = 1 \text{ translation} \quad (16)$$

$$l = 2 \text{ quadrupole deformation} \quad (17)$$

$$l = 3 \text{ octupole deformation} \quad (18)$$

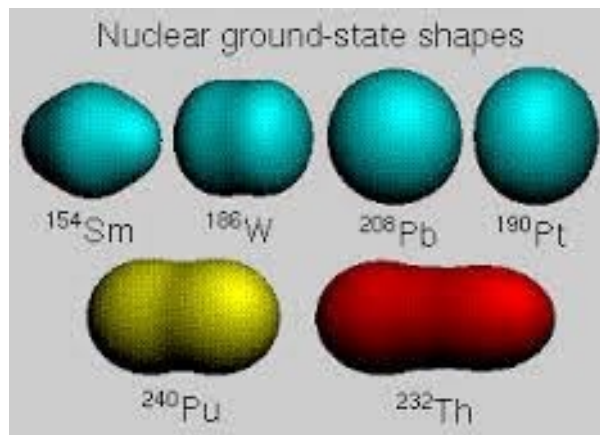


Figure 5:

The nuclear shapes can be determined from nuclear collective excitation, i.e., rotation and vibration.

4.4 Nuclear collective excitation

4.4.1 Rotation

4.4.2 Vibration